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Characters of the Positive Energy UIRs of
 $D=4$ Conformal Super-symmetry

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PLAN

Introduction to characters

General character formulae

Specifics for $su(2,2|N)$

Decompositions

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Elementary introduction to characters

Take the algebra $\mathfrak{sl}(2)$ with generators X^\pm, H and relations:

$$[X^+, X^-] = H, \quad [H, X^\pm] = \pm 2X^\pm$$

Lowest weight modules are defined by a lowest weight $\lambda \in \mathfrak{h}^*$ (\mathfrak{h} spanned by H) and lowest weight vector v_0 so that

$$Hv_0 = \lambda(H)v_0$$

$$X^-v_0 = 0$$

Lowest weight Verma modules V^λ are induced from 1-dim. reps of the Borel subalgebra \mathfrak{b} (spanned by H and X^-). In short

$$V^\lambda \cong U(\mathfrak{g}^+) \otimes v_0, \quad \mathfrak{g}^+ \text{ spanned by } X^+$$

Simple: it consists of vectors (omitting \otimes)

$$(X^+)^k v_0, \quad k = 0, 1, \dots$$

Take now a formal variable t . We can represent the Verma module by the infinite series:

$$1 + t + t^2 + \dots = \frac{1}{1-t}$$

So we define the character of the Verma module V^λ to be

$$\text{ch } V^\lambda \equiv \frac{e(\lambda)}{1-t}, \quad t \equiv e(\alpha)$$

Actually, $e(\cdot)$ is a formal exponent on \mathbb{Z}^k such that $e(0) = 1$, $e(\alpha)e(\alpha') = e(\alpha + \alpha')$
(though here we have only one root and also 1 is a multiple of α)

↓
omit

How to obtain f.d. reps from Verma modules? (2)

→ From reducible Verma modules, this happens

whenever holds the BGG condition

$$(\rho - \lambda, \beta) = \frac{m}{2} (\rho, \beta) \quad , \quad m \in \mathbb{N}, \quad \beta \in \Delta^+$$

$$\rho = \frac{1}{2} \sum_{\beta \in \Delta^+} \beta$$

in our case if $\lambda = \frac{m+1}{2} \beta$

Then V^λ contains a singular vector

$$v_3 = (X^+)^m v_0$$

$$\text{s.t.} \quad \begin{cases} H v_3 = \lambda'(H) v_3, & \lambda' = \lambda + m\beta \\ X^- v_3 = 0 \end{cases}$$

in other language V^λ contains a submodule

$$V^\lambda \supset V^{\lambda+m\beta}$$

In order to obtain an irred. representation we need to take the factor module

$$L_\lambda = V^\lambda / I^\lambda, \quad I^\lambda \text{ is the image of } V^{\lambda+m\beta}$$

i.e. L_λ will consist of m vectors

$$X^k v_0, \quad k = 0, 1, \dots, m-1$$

Obviously, the character of L_λ should (up to multiple)

$$1 + t + \dots + t^{m-1}$$

more logically, since L_λ is a factor-module we should have

$$\begin{aligned} \text{ch } L_\lambda &= \text{ch } V^\lambda - \text{ch } V^{\lambda+m\beta} = \\ &= \frac{e(\lambda)}{1-t} - \frac{e(\lambda+m\beta)}{1-t} = \\ &= \frac{e(\lambda)(1-e(m\beta))}{1-t} = \frac{e(\lambda)(1-t^m)}{1-t} \\ &= e(\lambda)(1+t+\dots+t^{m-1}) \end{aligned}$$

Generalization: Let \mathfrak{g} any simple Lie algebra with root system Δ , positive roots Δ^+

The character of a Verma module V^λ is

$$\text{ch } V^\lambda = e(\lambda) \prod_{\alpha \in \Delta^+} \frac{1}{1-e(\alpha)}$$

fundamental weights are obtained when BGG holds for every simple root (and consequently for any other positive root)

In this case the character formula can be written compactly as (5)

$$\text{ch } L_\Lambda = \sum_{w \in W} (-1)^{\ell(w)} \text{ch } V^{w \cdot \Lambda}$$

$w \in W \rightarrow$ Weyl group

This is what we had for $sl(2)$, where

The Weyl group consists of 2 elements: identity 1 and s_α , taking into account

$$\Lambda + w\beta = s_\beta \cdot \Lambda$$

$\ell(w)$: length function: $\sqrt{\text{minimal}} \#$ of simple reflections s_α which one can write w

Supercase \mathfrak{g} : basic classical superalgebra (excl. $osp(1/2n)$)

root system

$$\Delta^+ = \Delta_0^+ + \Delta_1^+ \quad \text{even, odd}$$

$$\text{ch } V^\Lambda = e(\Lambda) \left\{ \prod_{\alpha \in \Delta_0^+} \frac{1}{1 - e(\alpha)} \right\} \prod_{\alpha \in \Delta_1^+} (1 + e(\alpha))$$

explanation: odd elements are nilpotent, so only up to linear in all of them.

shorter: $\text{ch } V^\Lambda = \text{ch } V_0^\Lambda \text{ch } \hat{V}^\Lambda$

Suppose V^Λ is not reducible w.r.t. any odd root. (6)
Then the character formula is:

$$\begin{aligned} \text{ch } L_\Lambda &= \text{ch } L_\Lambda^0 \text{ch } \hat{V}^\Lambda = \\ &= \text{ch } L_\Lambda^0 \text{ch } \hat{L}_\Lambda \end{aligned}$$

Suppose now V^Λ is reducible w.r.t. odd root β ,
i.e. holds Kac condition

$$(\rho - \Lambda, \beta) = 0, \quad (\beta, \beta) = 0$$

Then we have embedding

$$V^\Lambda \supset V^{\Lambda+\beta}$$

In analogy with the even case we define

odd reflection [D + V. Petkova, 1985]

$$\hat{s}_\beta \cdot \Lambda = \Lambda + \beta, \quad (\beta, \beta) = 0, \quad (\Lambda, \beta) \neq 0$$

Indeed the odd embedding above is provided by a singular vector

$$v_{\beta}^{\beta} = P^{\beta} v_0 \quad \text{such that} \quad (P^{\beta})^2 = 0$$

Factorization of the submodule I^{β} generated by v_{β}^{β} (I^{β} is smaller than $V^{\beta+\beta}$!) means that we eliminate the root vector X_{β}^{+} from the basis, i.e. the odd character part will be

$$\text{ch } \hat{L}_{\Lambda} = \prod_{\substack{\alpha \in \Delta^{+} \\ \alpha \neq \beta}} (1 + e(\alpha))$$

This can be also written using the odd reflections:

using the factorization $\hat{L}_\lambda = \hat{V}^\lambda / \hat{I}^\lambda$

we have $ch \hat{L}_\lambda = ch \hat{V}^\lambda - ch \hat{I}^\lambda$

We know that \hat{I}^λ is "smaller" than $\hat{V}^{\lambda+\beta}$ - in characters it means:

WZ

$$ch \hat{I}^\lambda = \frac{1}{1+e(\beta)} ch \hat{V}^{\lambda+\beta} = \frac{e(\beta)}{1+e(\beta)} ch \hat{V}^\lambda$$

$$\Rightarrow ch \hat{L}_\lambda = ch \hat{V}^\lambda - \frac{1}{1+e(\beta)} ch \hat{V}^{\hat{\beta} \cdot \lambda}$$

$$= ch \hat{V}^\lambda - \hat{\beta} \cdot ch \hat{V}^\lambda =$$

$$= \sum_{\beta \in \hat{W}_\beta} (-1)^{l(\beta)} \beta \cdot ch \hat{V}^\lambda$$

$\hat{W}_\beta = \{1, \hat{\beta}\}$ since higher terms give zero

Thus, we have an easy conjecture for the general character formula:

$$ch \hat{L}_\lambda = \sum_{\beta \in \hat{W}} (-1)^{l(\beta)} \beta \cdot ch \hat{V}^\lambda$$

where $\hat{W} = \prod_{\beta \in K} \hat{W}_\beta$

↳ the set for which holds Kac criterion for reducibility

But life is more complicated!

There are 2 obstacles:

- 1) descendant redundancies, i.e., elimination of odd roots for which Kac condition does not hold, and even worse the opposite - non-elimination of such for which it holds!
- 2) redundancies, coming from the fact that we are looking for characters of a real form, and all states have to obey the condition of being physical.

Thus the general character formula may be written as:

$$\text{ch } \hat{L}_\Lambda = \sum_{\hat{\beta} \in \hat{W}} (-1)^{\ell(\hat{\beta})} \hat{\beta} \cdot (\text{ch } \hat{V}^\Lambda - \mathcal{R})$$

↘ counter-term

where $\hat{W} = \prod_{\beta \in K'} \hat{W}_\beta$

K' is the set of eliminated elements of the odd root lattice, i.e. roots and some non-roots (examples to be shown)

\mathcal{R} represents elements which has to be taken out since they are not allowed also in the odd Verma module

→ Important: even without any reducibility, \mathcal{R} may be non-zero!

Thus, we have to be more specific.

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Our algebra is the Lie superalgebra

$$\mathfrak{g} = \mathfrak{su}(2,2/N) = \mathfrak{g}_0 \oplus \mathfrak{g}_1 \quad (\text{even} + \text{odd})$$

with

$$\mathfrak{g}_0 = \mathfrak{su}(2,2) \oplus \mathfrak{u}(1) \oplus \mathfrak{su}(N)$$

The physically relevant representations are labelled by the signatures

$$\chi = [d; j_1, j_2; z; r_1, \dots, r_{N-1}]$$

where

d : conformal weight

j_1, j_2 : non-negative (half) integers labelling Lorentz f.d. irreps

z : represents $\mathfrak{u}(1)$

r_1, \dots, r_{N-1} : non-negative integers, Dynkin labels of $\mathfrak{su}(N)$

There is rich interplay between \mathfrak{g} and $\mathfrak{g}^{\mathbb{C}} = \mathfrak{sl}(4/N)$

Root system of $sl(4/N)$

$$\alpha_{ij}, \quad 1 \leq i < j \leq N+4$$

$$\text{even roots} \begin{cases} 1 \leq i < j \leq 4 & (\text{the 6 roots of } sl(4)) \\ 5 \leq i < j \leq 4+N & \text{the } \frac{N(N-1)}{2} \text{ roots of } sl(N) \end{cases}$$

$$4N \text{ odd roots} \quad 1 \leq i \leq 4, \quad 5 \leq j \leq 4+N$$

Reducibility of V^λ , $\lambda = \lambda(\alpha)$

$sl(4)$: we calculate $n_{ij} = (\rho - \lambda, \alpha_{ij}^\vee)$, $\alpha_{ij}^\vee = \frac{2\alpha_{ij}}{(\alpha_{ij}, \alpha_{ij})}$
(34(2.1))

$$\text{simple roots} \begin{cases} n_{12} = 1 + 2j_1 \equiv n_1 \in \mathbb{N} \\ n_{23} = 1 - d - j_1 - j_2 \equiv n_2 \\ n_{34} = 1 + 2j_2 \equiv n_3 \in \mathbb{N} \end{cases}$$

$$n_{13} = 2 - d + j_1 - j_2 = n_1 + n_2$$

$$n_{24} = 2 - d - j_1 + j_2 = n_2 + n_3$$

$$n_{14} = 3 - d + j_1 + j_2 = n_1 + n_2 + n_3$$

f.d. irreps \equiv all $n_i \in \mathbb{N}$, but in all our cases we have:

$n_1, n_3 \in \mathbb{N}, \quad n_2 \notin \mathbb{N}$

omit

Odd reducibilities

$4N$ roots, $4N$ possibilities, 4 series:

$$d = d_{NK}^1 - z \delta_{NY} \quad K=1, \dots, N$$

$$d_{NK}^1 \equiv 4 - 2r + 2j_2 + r + 2m_K - 2m/N$$

$$d = d_{NK}^2 - z \delta_{NY}$$

$$d_{NK}^2 \equiv 2 - 2K - 2j_2 + r + 2m_K - 2m/N$$

$$d = d_{NK}^3 - z \delta_{NY}$$

$$d_{NK}^3 \equiv 2 + 2K - 2N + 2j_2 - r - 2m_K + 2m/N$$

$$d = d_{NK}^4 - z \delta_{NY}$$

$$d_{NK}^4 \equiv 2K - 2N - 2j_2 - r - 2m_K + 2m/N$$

where

$$m_K \equiv \sum_{i=1}^{N-1} r_i$$

number of cells of the K -th row of the standard Young tableau

$$m \equiv \sum_{K=1}^{N-1} m_K = \sum_{K=1}^{N-1} K r_K$$

number of cells

Note that for fixed module only one condition from each series can be satisfied. Thus at most

4 odd conditions may be satisfied.

Fortunately, for the UIR we shall require only coincidence of up to 2 odd reducibilities.

Unitarity

Theorem [V.O. + V. Petkova, 1985]

The complete list of lowest weight UIRs of $su(2, 2|N)$ is:

$$d \geq d_{\max} = \max(d_{N1}^1, d_{NN}^3)$$

$$d = d_{NN}^4 \geq d_{N1}^1, \quad j_1 = 0$$

$$d = d_{N1}^2 \geq d_{NN}^3, \quad j_2 = 0$$

$$d = d_{N1}^2 = d_{NN}^4, \quad j_1 = j_2 = 0$$

where d_{\max} is the threshold of the continuous unitary spectrum.

Note: in the last case we have

$$d = m_1, \quad t = \frac{2m_1}{N} - m_1$$

and for $N=1$ by agreement $m_1 = m = 0$ (no Tableau for $su(1)$!), i.e. all quantum numbers are missing zeros, i.e. this is the trivial 1-dim. UIR.

For $N=1$ the above list was given first by Flato and Fronsdal, 1984.

Even for $d > d_{\max}$ things are not trivial for low values of spin and isospin, but the most interesting things happen at reducible points.

Thus, only 4 odd reducib. points are playing a role, and they have partial orders

$$d_{N1}^1 > d_{N1}^2, \quad d_{NN}^3 > d_{NN}^4$$

Thus, at most 2 such points may coincide. Thus, we have single reducibility conditions (SRC) and double-reducibility-conditions (DRC) situations

Take first SRC cases (4 of them)

$d = d_{max} = d_{N1}^1 > d_{NN}^3$	} $\alpha_{3,4+N} = \delta_4 + \delta_2$ ^{odd loc.}	
$d = d_{N1}^2 > d_{NN}^2, \quad \tilde{j}_2 = 0$		$\alpha_{4,4+N} = \delta_4$
$d = d_{max} = d_{NN}^2 > d_{N1}^1$		$\alpha_{15} = \delta_3 + \delta_1$ _{even loc.}
$d = d_{NN}^4 > d_{N1}^1, \quad \tilde{j}_1 = 0$		$\alpha_{25} = \delta_3$

Even singular vectors

$$v_j^1 = (X_1^+)^{1+2j_1} v_0$$

$$v_j^2 = (X_3^+)^{1+2j_2} v_0$$

$$v_j^j = (X_j^+)^{1+\Gamma_{N+4-j}} v_0 \quad j=5, \dots, N+3$$

odd singular vectors

$$v_{\text{odd}}^1 = P_{3,4+N} v_0 = (X_4^+ X_2^+ (h_2 - 1) - X_2^+ X_4^+ h_2) v_0$$

Cartan elem.

$$v_{\text{odd}}^2 = X_4^+ v_0$$

$$v_{\text{odd}}^3 = P_{15} v_0 = (X_3^+ X_1^+ (h_1 - 1) - X_1^+ X_3^+ h_1) v_0$$

$$v_{\text{odd}}^4 = X_3^+ v_0$$

exercise: check that polynomials are nilpotent
(trivial for ex. 2, 4)

We now introduce the factor module L_1 .

Let us denote by $|0\rangle$ the lowest weight vector of L_1 . Each singular vector of V^1 acts, generates zero action on $|0\rangle$ (mult)

Even

$$(X_1^+)^{1+2j_1} |0\rangle = 0$$

$$(X_2^+)^{1+2j_2} |0\rangle = 0$$

$$(X_j^+)^{1+\Gamma_{N+4-j}} |0\rangle = 0, \quad j=5, \dots, N+3$$

Odd
whenever
applicable

$$P_{3,4+N} |0\rangle = (2j_2 X_{3+4+N}^+ - X_4^+ X_2^+) |0\rangle = 0$$

$$X_4^+ |0\rangle = 0$$

$$P_{1,5} |0\rangle = (2j_1 X_{1+5}^+ - X_5^+ X_1^+) |0\rangle = 0$$

$$X_5^+ |0\rangle = 0$$

We see that whenever the appropriate condition holds, then the corresponding odd root vector is annihilated, i.e.

$$X_{3+4+N}^+ \quad \text{when } j_2 \neq 0 \quad !$$

$$X_4^+$$

$$X_{1+5}^+ \quad \text{when } j_1 \neq 0 \quad !$$

$$X_5^+$$

Unexpected problems Explanation!

Thus, in these 2 cases, singular vectors and null action are descendant, but they are replaced by more subtle null action, not generated by the kac root conditions =

$$X_3^+ X_1^+ X_3^+ |0\rangle = 0, \quad d = d_{NN}^3, \quad \tilde{J}_1 = 0$$

$$X_4^+ X_2^+ X_4^+ |0\rangle = 0, \quad d = d_{NN}^1, \quad \tilde{J}_2 = 0$$

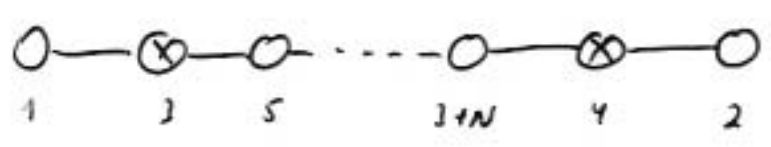
This is a peculiarity of the super case.

Thus, in these latter cases the character formula will be (the odd part of it)

$$ch \hat{L}_\Lambda = \sum_{w \in W'} (-1)^w \hat{S}_w \cdot ch \hat{V}_\Lambda^A \quad (r_i \neq 0)$$

$$W' = \left\{ 1, \beta_{\alpha_1 + 2\alpha_2} \right\} \text{ or } \left\{ 1, \beta_{\alpha_2 + 2\alpha_1} \right\} \quad \text{not roots !!}$$

More complications if some $r_i = 0$. Here we need explicitly the Dynkin diagram:



This is a non-distinguished D.d. with 2 odd simple roots, most suitable for our purposes.

Example, $r_1 = 0$, $r_2 \neq 0$
 Consider the vector

$$P_{3,4+N} V_0 = (2j_2 X_{3,3+N} - X_{4,3+N} X_6^+) V_0 = \\ = (P_{3,4+N} X_{3+N}^+ - X_{3+N}^+ P_{3,4+N}) V_0$$

and we consider the case when $P_{3,4+N} |0\rangle = 0$
 and since $X_{3+N}^+ |0\rangle = 0$ ($\Leftarrow r_1 = 0$)

$$\text{Then } \Rightarrow P_{3,3+N} |0\rangle = 0$$

i.e. for $r_2 \neq 0$ we eliminate also the
 root vector $X_{3,3+N}$

Analogously, if

$$r_1 = r_2 = \dots = r_{i_0} = 0$$

then all vectors $X_{3,4+N-i}$, $i \leq i_0$

are eliminated; in particular,

if all $r_i = 0$, then all $X_{3,4+k}$ $k=1, \dots, N$
 are eliminated.

The character f - h is

$$\text{ch } \hat{L}_M = \sum_{\beta \in \hat{W}_{i_0}} (-1)^{l(\beta)} \beta \cdot \text{ch } \hat{V}^{\beta}$$

$$\hat{W}_{i_0}^a = \hat{W}_{\alpha_{3,4+k}} \times W_{\alpha_{3,4+k}} \times \dots \times \hat{W}_{\alpha_{3,4+k-i_0}}$$

Thus, we have illustrated peculiarities of the 1st kind, namely, elimination of vectors, that do not follow from Kac red. cond.

Now we turn to problems of second kind, namely, those that originate from the fact that we consider $SU(2,2/N)$, not present for $sl(4/N)$.

This can be illustrated already for $N=1$. In that case we have 7 odd vectors. The claim is that sometimes they can not be applied on the vacuum to produce a component of a superfield.

The vacuum $|0\rangle$ carries the quantum numbers of the lowest weight vector v_0 i.e. for $N=1$

$$X = [d; j_1, j_2; z] \xrightarrow{\beta} X'$$

$$X_{d_{3,N+4}}^+ : X' = [d + \frac{1}{2}; j_1, \underline{j_2 - \frac{1}{2}}; z + \frac{3}{2}]$$

$$X_{d_{4,N+4}}^+ : X' = [d + \frac{1}{2}; j_1, j_2 + \frac{1}{2}; z + \frac{3}{2}]$$

$$X_{d_{1,5}}^+ : X' = [d + \frac{1}{2}; \underline{j_1 - \frac{1}{2}}, j_2; z - \frac{3}{2}]$$

$$X_{d_{2,5}}^+ : X' = [d + \frac{1}{2}; j_1 + \frac{1}{2}, j_2; z - \frac{1}{2}]$$

Obviously, if $j_2 = 0$, the state

$$X_{2, N+1}^+ |0\rangle \text{ is not allowed}$$

and if $j_1 = 0$, the state

$$X_{1, 1}^+ |0\rangle \text{ is not allowed}$$

For $N > 1$ there are similar restrictions from the values of r_i . Altogether, the general odd state looks as follows

$$\Psi = \left(\prod_{k=N}^1 (X_{1, 4+k}^+)^{\epsilon_{1, 4+k}} \right) \left(\prod_{k=N}^1 (X_{2, 4+k}^+)^{\epsilon_{2, 4+k}} \right) \gamma$$

$$\times \left(\prod_{k=1}^N (X_{3, 4+k}^+)^{\epsilon_{3, 4+k}} \right) \left(\prod_{k=1}^N (X_{4, 4+k}^+)^{\epsilon_{4, 4+k}} \right) |0\rangle$$

where $\epsilon_{i,j} = 0, 1$ indicates the presence of a gen.
Conditions for a state to be allowed:

$$\sum_{k=1}^N \epsilon_{1, 4+k} \leq \sum_{k=1}^N \epsilon_{2, 4+k} + 2j_1 \tag{*}$$

$$\sum_{k=1}^N \epsilon_{3, 4+k} \leq \sum_{k=1}^N \epsilon_{4, 4+k} + 2j_2$$

$$\epsilon_{1, 5+i} + \epsilon_{2, 5+i} + \epsilon_{3, 5+i} + \epsilon_{4, 5+i} \leq \epsilon_{1, 4+i} + \epsilon_{2, 4+i} + \epsilon_{3, 5+i} + \epsilon_{4, 5+i} + r_i$$

$i = 1, \dots, N-1$

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These conditions are also sufficient only for $N=1$.
To formulate the exact condition we need
the notion of level ε

$$\varepsilon = \sum_{k=1}^N (\varepsilon_{1,4+k} + \varepsilon_{2,4+k} + \varepsilon_{3,4+k} + \varepsilon_{4,4+k})$$

i.e. the total number of all generators applied
on $|0\rangle$

Criterion: The necessary and sufficient conditions
for a state of level ε to be allowed
are that conditions (*) are fulfilled and that
the state is a descendant of an allowed
state of level $\varepsilon-1$.

Examples of states which do not fulfil
the second part of the Criterion only, are
given in the text.

Thus, even for UIRs which come from oddly
irreducible Verma modules there will be lots
of forbidden states at low spins and isospin

The maximal number of states is 2^{4N}
and it is realized only if

$$j_1, j_2 \geq \frac{N}{2}, \quad r_i \geq 4, \quad \forall i$$

Decompositions

SRC:

$$\hat{L}_{\text{long}} = \hat{L}_\Lambda \oplus \hat{L}_{\Lambda+\beta}$$

↘ as vector spaces

always for $N=1$
for $N>1$: some
conditions

$$d = d_{\text{max}} = d^1 > d^2, \quad \Gamma_1 > 0$$

$$j_2 > 0, \quad \beta = \alpha_{2, 4+N} \quad (\text{the root corresp. to } d^1)$$

$$j_2 = 0, \quad \beta = \alpha_{2, 4+N} + \alpha_{4, 4+N} \quad (\leftarrow \text{not a root!})$$

$$d = d_{\text{max}} = d^2 > d^1, \quad \Gamma_{N-1} > 0$$

$$j_1 > 0, \quad \beta = \alpha_{1, 5} \quad (\text{the root corresp. to } d^2)$$

$$j_1 = 0, \quad \beta = \alpha_{1, 5} + \alpha_{2, 5} \quad (\leftarrow \text{not a root!})$$

DKC:

$$\hat{L}_{\text{long}} = \hat{L}_\Lambda + \hat{L}_{\Lambda+\beta} + \hat{L}_{\Lambda+\beta'} + \hat{L}_{\Lambda+\beta+\beta'}$$

$$\Gamma_1 \Gamma_{N-1} > 0$$

$$d = d_{\text{max}} = d^1 = d^2$$

β, β' : 4 cases combinations of the above

UIRs that do not enter decompositions

SRC cases, $N > 1$

a) $d = d_{\max} = d^1 > d^2$
 $i_1, i_2 \sim, \quad \Gamma_1 = 0$

b) $d = d^2 > d^1$
 $i_1 \sim, i_2 = 0, \quad \Gamma_1 \leq 2$

c) $d = d_{\max} = d^2 > d^1$
 $i_1, i_2 \sim, \quad \Gamma_{N-1} = 0$

d) $d = d^1 > d^2$
 $i_1 = 0, i_2 \sim, \quad \Gamma_{N-1} \leq 2$

DRC cases

all cases when $N=1$

for $N > 1$ the list is:

$$a) \quad d = d_{\max} = d^1 = d^3 = 2 + j_1 + j_2 + m_1$$

$$j_1, j_2 \sim, \quad \Gamma_1 \Gamma_{N-1} = 0$$

$$ad) \quad d = d^1 = d^4 = 1 + j_2 + m_1$$

$$j_1 = 0, \quad j_2 \sim, \quad \Gamma_{N-1} \leq 2 \quad \forall N > 1$$

$$\Gamma_1 = 0, \quad N=2$$

$$bc) \quad d = d^2 = d^3 = 1 + j_1 + m_1$$

$$j_1 \sim, \quad j_2 = 0, \quad \Gamma_1 \leq 2 \quad \forall N > 1$$

$$\Gamma_{N-1} = 0, \quad N=2$$

$$bd) \quad d = d^2 = d^4 = m_1$$

$$j_1 = j_2 = 0,$$

$$\Gamma_1 \leq 4, \quad N=2$$

$$\Gamma_1, \Gamma_{N-1} \leq 2, \quad N > 2$$